

Evolution problems

Rob Stevenson
Korteweg-de Vries Institute for Mathematics
University of Amsterdam

Overview

- With time marching methods, an optimal distribution of grid points over space and time is hard to realize. ■
- We apply an adaptive method to a simultaneously space-time variational formulation. ■
- While keeping discrete solutions on all time levels is prohibitive for time marching methods, thanks to the use of tensorized multi-level bases our method produces approximations simultaneously in space and time without penalty in complexity because of the additional time dimension.

Parabolic problems

Let $V \hookrightarrow H \simeq H' \hookrightarrow V'$, $I := (0, T)$. Consider parabolic problem

$$\begin{cases} \frac{\partial u}{\partial t}(t, \cdot) + A(t)u(t, \cdot) = g(t, \cdot) & \text{in } V' \quad (t \in I) \\ u(0, \cdot) = u_0 & \text{in } H, \end{cases}$$

where for $t \in I$ a.e.,

$$|(A(t)(\eta))(\zeta)| \leq M_a \|\eta\|_V \|\zeta\|_V \quad (\eta, \zeta \in V) \quad (\text{boundedness}),$$

$$(A(t)(\eta))(\eta) + \mu \|\eta\|_H^2 \geq \alpha \|\eta\|_V^2 \quad (\eta \in V) \quad (\text{Gårding inequality}).$$

We take $H = L_2(\Omega)$, and think of $A(t)$ differential or integrodifferential operator of order $2m \geq 0$, and $V = H^m(\Omega)$ ($H_0^m(\Omega)$).

Ex 1. $m = 1$, $(A(t)(\eta))(\zeta) = \int_{\Omega} \nabla \eta \cdot \nabla \zeta dx$ (heat eq).

Space-time variational formulation

Multiplying with smooth test functions that vanish at $t = T$, and integration over space and time, and int.-by-parts in time (initial condition becomes a natural b.c.),

$$\begin{aligned}(Bu)(v) &:= - \int_I \int_{\Omega} u(t, x) \frac{\partial v}{\partial t}(t, x) dx dt + \int_I (A(t)u(t, \cdot))(v(t, \cdot)) dt \\ &= \int_I \int_{\Omega} g(t, x)v(t, x) dx dt + \int_{\Omega} u_0(x)v(0, x) dx =: f(v) \blacksquare\end{aligned}$$

With

$$\mathcal{X} := L_2(I; V) = \left\{ u : I \rightarrow V : \int_I \|u(t, \cdot)\|_V^2 dt < \infty \right\} \simeq L_2(I) \otimes V,$$

$$\mathcal{Y} := L_2(I; V) \cap H_{0, \{T\}}^1(I; V')$$

one has $B \in \mathcal{B}(\mathcal{X}, \mathcal{Y}')$. \blacksquare Also

Thm 1. $B^{-1} \in \mathcal{B}(\mathcal{Y}', \mathcal{X})$.

Proof. (for $A(t) \equiv A = A'$ coercive, and $V \hookrightarrow L_2(\Omega)$ compact)

\exists orthon. basis $\{\phi\}$ for $L_2(\Omega)$ of eigenfunctions of A with eigenv. $\lambda_\phi > 0$. ■

Writing $u(t, x) = \sum_\varphi u_\varphi(t)\varphi(x)$, $f(t, x) = \sum_\varphi f_\varphi(t)\varphi(x)$, $v(t, x) = \sum_\varphi v_\varphi(t)\varphi(x)$,

$$(Bu)(v) = f(v) \Leftrightarrow \int_I -u_\varphi \dot{v}_\varphi + \lambda_\varphi u_\varphi v_\varphi dt = \int_I f_\varphi v_\varphi dt \quad (\varphi \in \{\varphi\}, v_\varphi \in L_2(I)). \blacksquare$$

Equipping V by $(A\cdot)(\cdot)^{\frac{1}{2}}$ and $H_{0,\{T\}}^1(I)$ by $|\cdot|_{H^1(I)}$,

$$\begin{aligned} \|u\|_{L_2(I;V)}^2 &= \sum_\varphi \lambda_\varphi \|u_\varphi\|_{L_2(I)}^2 \\ \|v\|_{L_2(I;V) \cap H_{0,\{T\}}^1(I;V')}^2 &= \sum_\varphi \lambda_\varphi \|v_\varphi\|_{L_2(I)}^2 + \lambda_\varphi^{-1} |v_\varphi|_{H^1(I)}^2 \end{aligned}$$

Now sufficient to prove that

with

$$(B_\lambda u)(v) := \int_I -u\dot{v} + \lambda uv \, dt,$$

$B_\lambda : L_2(I) \rightarrow (H_{0,\{T\}}^1(I), \|\cdot\|_\lambda)'$ is b.i., uniformly in $\lambda > 0$, where

$$\|\cdot\|_\lambda := \sqrt{\lambda^2 \|\cdot\|_{L_2(\Omega)}^2 + \|\cdot\|_{H^1(I)}^2}. \quad \blacksquare$$

Invertibility of B_λ , and its unif. boundedness are readily shown.

Given $v \in H_{0,\{T\}}^1(I)$, define $u = -\dot{v} + \lambda v$. Then $\|u\|_{L_2(I)} \leq \sqrt{2} \|v\|_\lambda$, and

$$(B_\lambda u)(v) = \int_I \dot{v}(t)^2 - 2\lambda v(t)\dot{v}(t) + \lambda^2 v(t)^2 dt \geq \|v\|_\lambda^2 \geq \frac{1}{2}\sqrt{2} \|v\|_\lambda \|u\|_{L_2(I)},$$

where we used $\int_I -2v(t)\dot{v}(t) dt = -\int_I \frac{d}{dt} v(t)^2 dt = -v(T)^2 + v(0)^2 = v(0)^2 \geq 0$.

This unif. inf-sup condition implies the unif. boundedness of $B_\lambda^{-1} : (H_{0,\{T\}}^1(I), \|\cdot\|_\lambda)' \rightarrow L_2(I)$. \square

Shift of the smoothness indices in space

Thm 2. Let $W \hookrightarrow V \hookrightarrow U$, $A(t)' + \mu I : W \rightarrow U'$ b.i., and $A(\cdot)' \in C(\bar{I}, \mathcal{B}(W, U'))$. Then with

$$\mathcal{X} := L_2(I; U')$$

$$\mathcal{Y} := L_2(I; W) \cap H_{0, \{T\}}^1(I; U)$$

$B \in \mathcal{B}(\mathcal{X}, \mathcal{Y}')$ is b.i.

For $W = V = U$, one retrieves original spaces.

Typical application: $m = 1$, $V = H_0^1(\Omega)$, $W = H^2(\Omega) \cap H_0^1(\Omega)$, $U = L_2(\Omega) \simeq U'$, assuming H^2 -regularity of the spatial operator.

Tensor product bases

Let $\Theta^{\mathcal{X}}$, $\Theta^{\mathcal{Y}}$, and $\Sigma^{\mathcal{X}}$, $\Sigma^{\mathcal{Y}}$ be collections of temporal or spatial functions such that, normalized in the corresponding norms,

$$\begin{array}{llll} \Theta^{\mathcal{X}} & \text{is a Riesz basis for } L_2(I), & & \\ \Theta^{\mathcal{Y}} & \text{"} & L_2(I) \text{ and for } H_{0,\{T\}}^1(I), & \\ \Sigma^{\mathcal{X}} & \text{"} & U', & \\ \Sigma^{\mathcal{Y}} & \text{"} & W & \text{" } U. \end{array}$$

Then, with $\mathcal{X} = L_2(I; U')$ and $\mathcal{Y} = L_2(I; W) \cap H_{0,\{T\}}^1(I; U)$, normalized in the corresponding norms,

$$\begin{array}{ll} \Theta^{\mathcal{X}} \otimes \Sigma^{\mathcal{X}} & \text{is a Riesz basis for } \mathcal{X}, \\ \Theta^{\mathcal{Y}} \otimes \Sigma^{\mathcal{Y}} & \text{" } L_2(I; W), H_{0,\{T\}}^1(I; U), \text{ and so for } \mathcal{Y}, \end{array}$$

Best possible rates

Let $U = L_2(\Omega) \simeq U'$, $\Omega \subset \mathbb{R}^n$, so $\mathcal{X} = L_2(I; L_2(\Omega))$

$$s_{\max} = \begin{cases} \min(d_t, \frac{d_x}{n}) & \text{isotropic spatial wavelets} \\ \min(d_t, d_x) & \text{anisotropic spatial wavelets} \end{cases}$$

up to log-factors when $d_t = \frac{d_x}{n}$ or $d_t = d_x$. ■

With non-adaptive approx, these rates require boundedness of certain mixed derivatives in L_2 . Relaxed regularity conditions with best N -term approx. ■

Realization of rates of best N -term approximation: Application of awgm to

$$\mathbf{B}^* \mathbf{B} \mathbf{u} = \mathbf{B}^* \mathbf{f}.$$

Numerics heat equation

Heat equation

$$\begin{aligned}\frac{\partial}{\partial t}u - \Delta_{\mathbf{x}}u &= g \quad \text{on } (0, T) \times \square, \\ u &= 0 \quad \text{on } (0, T) \times \partial\square \\ u(0, \cdot) &= u_0.\end{aligned}$$

$\square = (0, 1)^n$. Temporal wavelets of order $d_t = 5$, tensor product spatial wavelets of order $d_x = 5$.

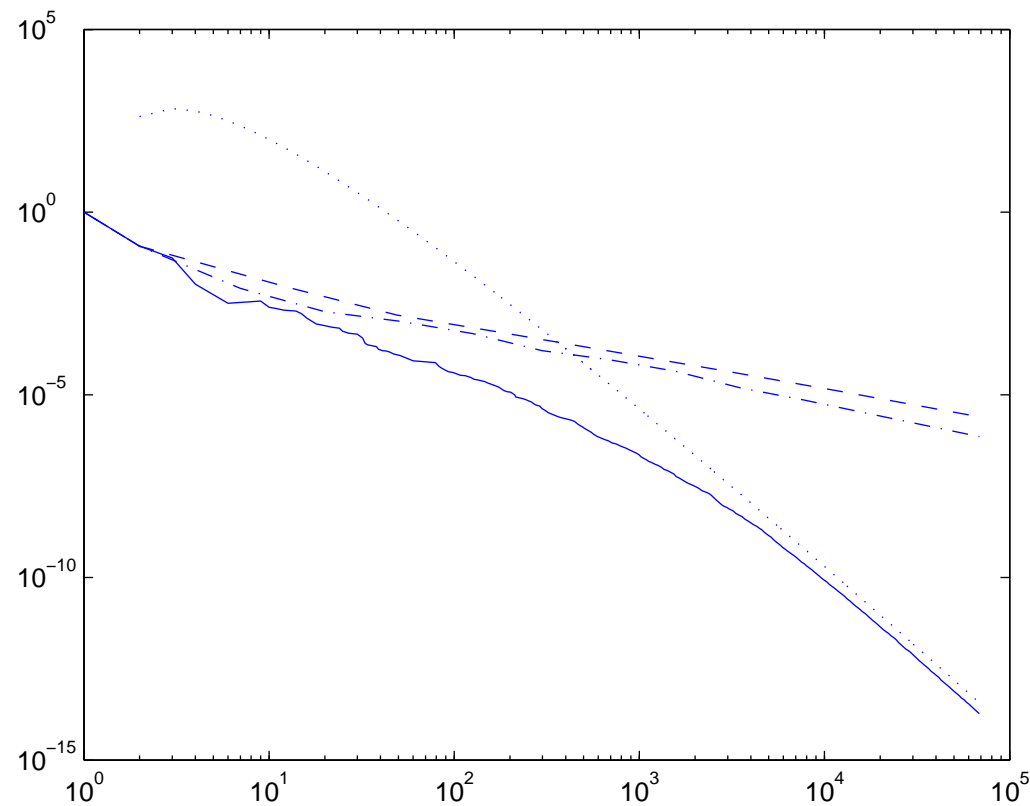


Figure 1: Heat eqn. in $n = 1$ spatial dimension, right-hand side $g = 1$ and initial condition $u_0 = 0$. $\|\mathbf{B}\mathbf{u}_\varepsilon - \mathbf{f}\| / \|\mathbf{f}\|$ vs. $N = \#\text{supp } \mathbf{u}_\varepsilon$ for awgm (solid), full-grid (dashed) and sparse-grid method (dashed-dotted). The dotted line is a multiple of $N^{-5}(\log N)^{5\frac{1}{2}}$.

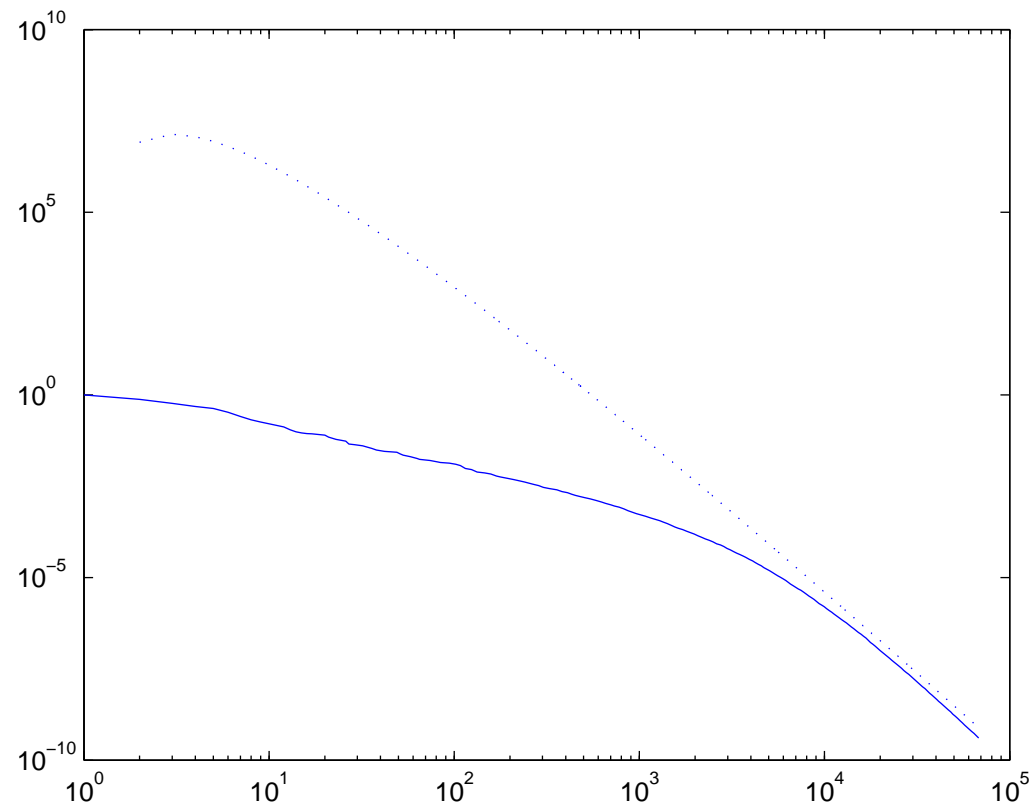


Figure 2: awgm applied to heat eqn. in $n = 1$ spatial dimension, right-hand side $g = 1$ and initial condition $u_0 = 1$. $\|\mathbf{B}\mathbf{u}_\varepsilon - \mathbf{f}\|/\|\mathbf{f}\|$ vs. $N = \#\text{supp } \mathbf{u}_\varepsilon$. The dotted line is a multiple of $N^{-5}(\log N)^{5\frac{1}{2}}$.

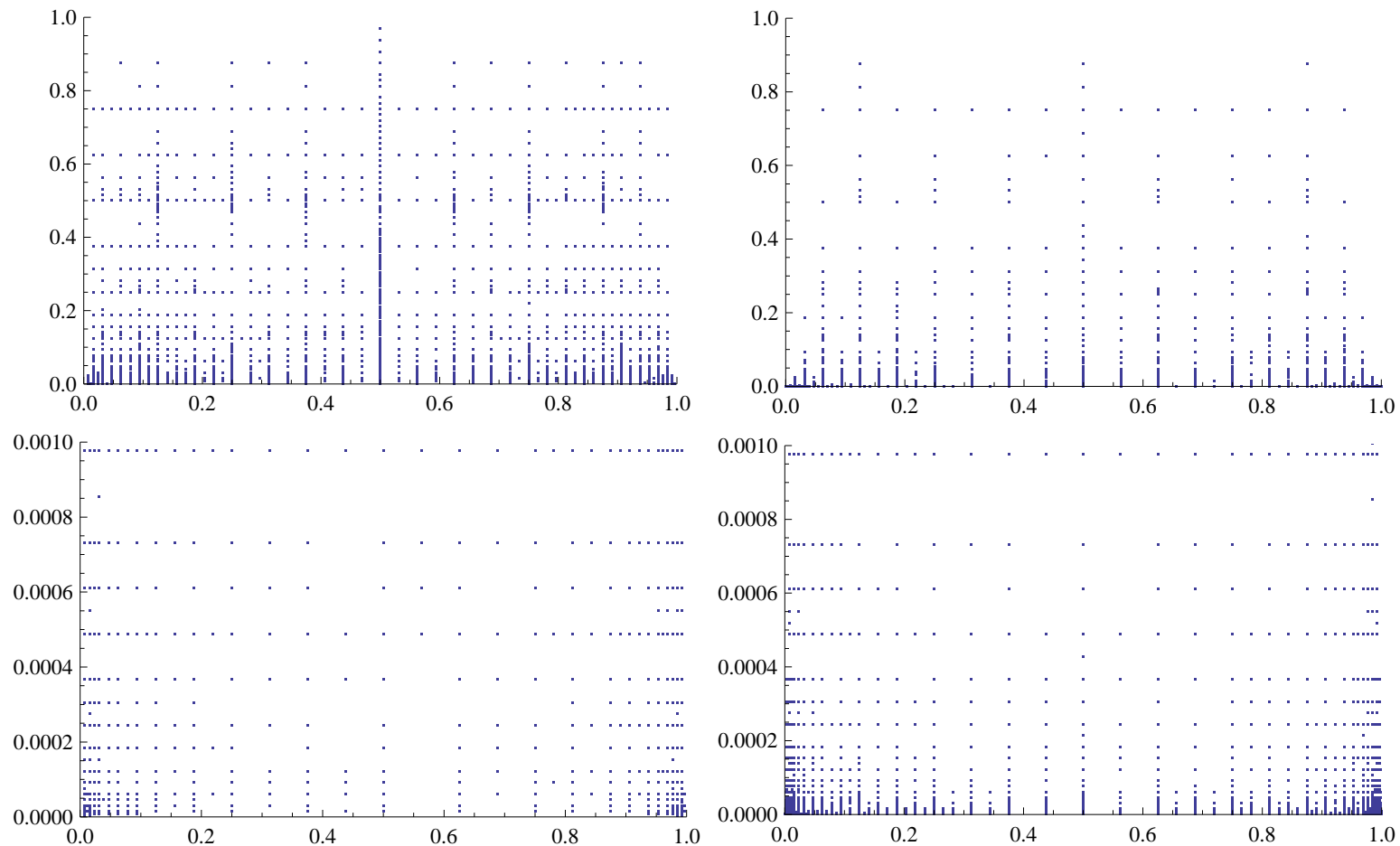


Figure 3: Heat eqn. in $n = 1$ spatial dimension and right-hand side $g = 1$. Centers of the supports of the wavelets selected by `awgm`. Left $u_0 = 0$ and $\#\mathbf{u}_\varepsilon = 13420$. Right $u_0 = 1$ and $\#\mathbf{u}_\varepsilon = 13917$. A zoom in near $t = 0$ is given at the bottom row.

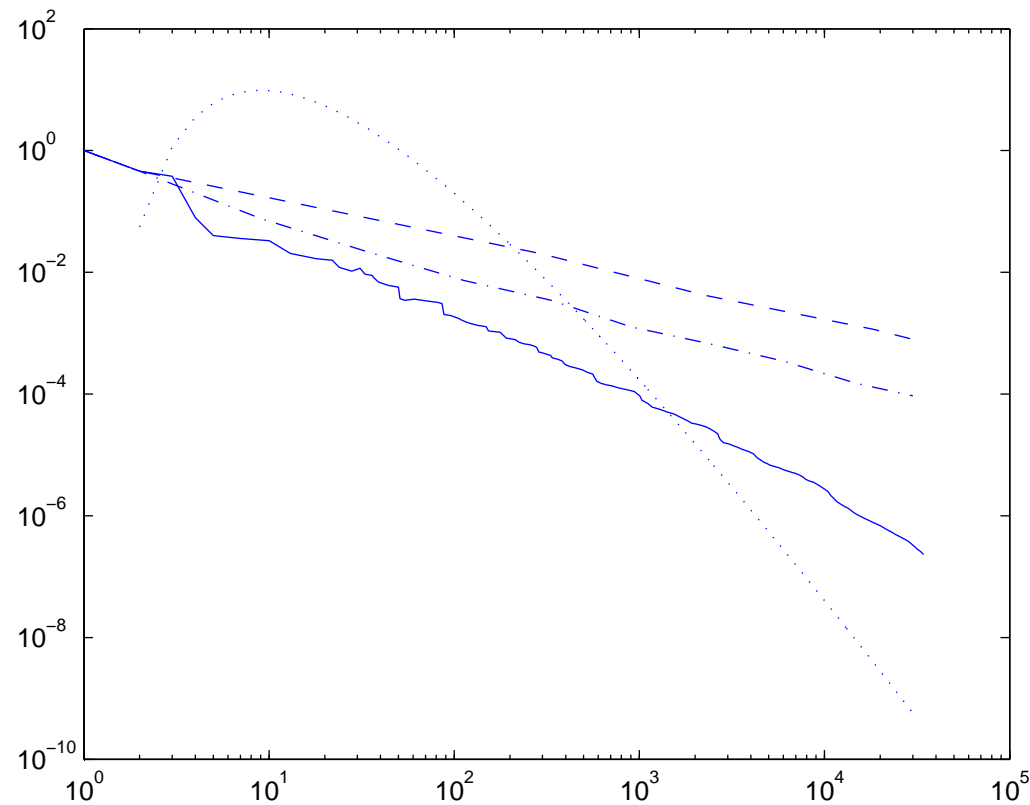


Figure 4: Heat eqn. in $n = 2$ spatial dimensions, right-hand side $g = 1$ and initial condition $u_0 = 0$. $\|\mathbf{B}\mathbf{u}_\varepsilon - \mathbf{f}\| / \|\mathbf{f}\|$ vs. $N = \#\text{supp } \mathbf{u}_\varepsilon$ for awgm (solid), full-grid (dashed) and sparse-grid method (dashed-dotted). The dotted line is a multiple of $N^{-5}(\log N)^{11}$.

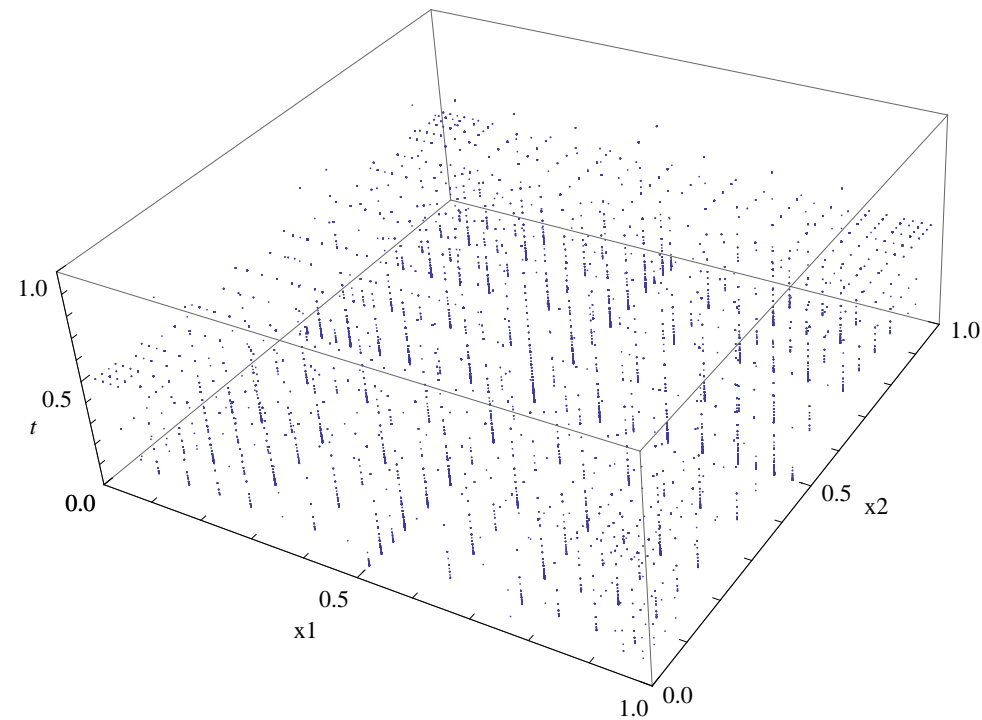


Figure 5: Heat eqn. in $n = 2$ spatial dimensions, right-hand side $g = 1$, and initial condition $u_0 = 0$. Centers of the supports of the wavelets selected by awgm. $\#\mathbf{u}_\varepsilon = 34316$.

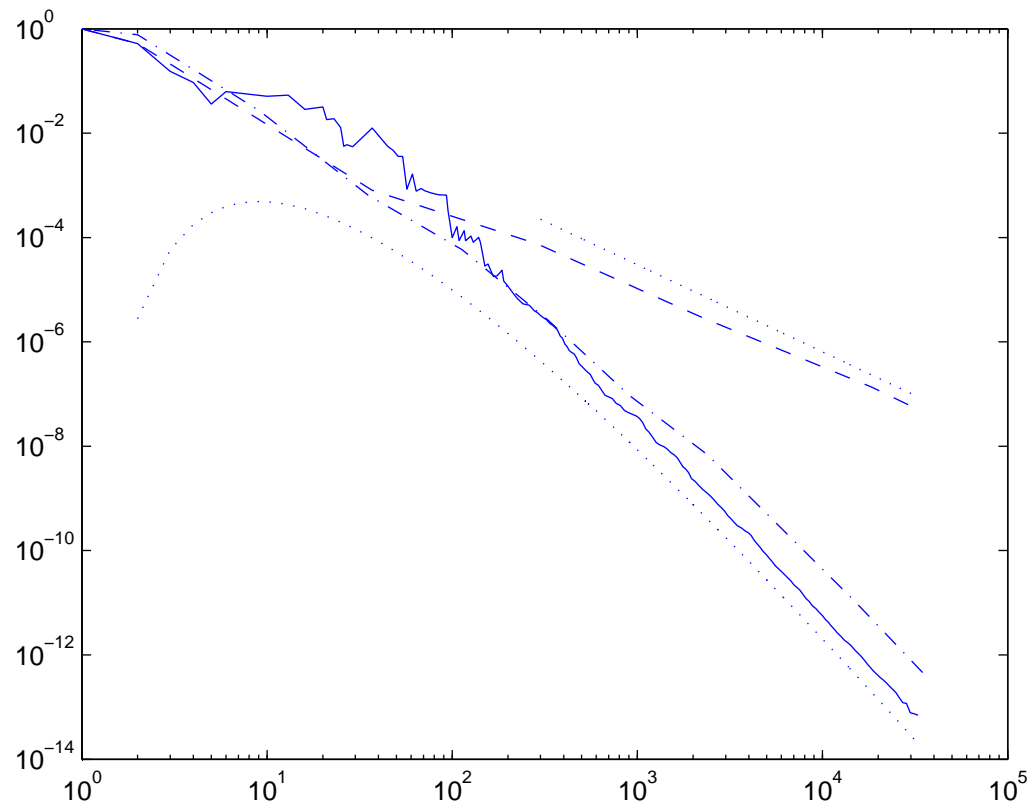


Figure 6: Heat eqn. in $n = 2$ spatial dimensions, right-hand side $g(t, \mathbf{x}) = t^4 x_1(x_1 - 1)(x_1^2 - x_1 - 1)x_2(x_2 - 1)(x_2^2 - x_2 - 1)$ and initial condition $u_0 = 0$. $\|\mathbf{B}\mathbf{u}_\varepsilon - \mathbf{f}\| / \|\mathbf{f}\|$ vs. $N = \#\text{supp } \mathbf{u}_\varepsilon$ for awgm (solid), full-grid (dashed) and sparse-grid method (dashed-dotted). The dotted lines are multiples of $N^{-5}(\log N)^{11}$ and $N^{-5/3}$, respectively.

Instationary Stokes equations

$$\left\{ \begin{array}{ll} \frac{\partial \vec{u}}{\partial t} - \nu \Delta \vec{u} - \nabla p & = \vec{g} \quad \text{on } I \times \Omega \\ \operatorname{div} \vec{u} & = 0 \quad \text{on } I \times \Omega \\ \vec{u} & = 0 \quad \text{on } I \times \partial\Omega \\ \vec{u}(0, \cdot) & = u_0 \quad \text{on } \Omega \end{array} \right.$$

Multiplying with smooth, *divergence-free* test functions v , integrating over $I \times \Omega$, int-by-parts over space and time, yields

$$\begin{aligned} (B\vec{u})(\vec{v}) &:= - \int_I \int_{\Omega} \vec{u} \cdot \frac{\partial \vec{v}}{\partial t} dx dt + \int_I \int_{\Omega} \nu \nabla \vec{u} : \nabla \vec{v} dx dt \\ &= \int_I \int_{\Omega} \vec{g} \cdot \vec{v} dx dt + \int_{\Omega} \vec{u}_0 \cdot \vec{v}(0, \cdot) dx =: \vec{f}(\vec{v}) \blacksquare \end{aligned}$$

Set $H := H_0(\operatorname{div} 0; \Omega) = \{ \vec{u} \in L_2(\Omega)^n : \operatorname{div} \vec{u} = 0 \text{ on } \Omega, \vec{u} \cdot \vec{n} = 0 \text{ on } \partial\Omega \}$,

$$V := H_0^1(\Omega)^n \cap H.$$

$(A\vec{\eta})(\vec{\zeta}) = \int_{\Omega} \nu \nabla \vec{\eta} : \nabla \vec{\zeta} dx$ satisfies the boundedness assump. and the Gårding inequality.

So with $\mathcal{X} = L_2(I; V)$ and $\mathcal{Y} = L_2(I; V) \cap H_{0, \{T\}}^1(I; V')$, $B : \mathcal{X} \rightarrow \mathcal{Y}$ b.i. ■

Problem: The construction of a wavelet Riesz basis for V . Moreover, method won't yield p .

Space-time variational saddle point formulation

Drop divergence-free condition on the test functions, and impose divergence-free condition on \vec{u} weakly.

$$\begin{aligned} (B(\vec{u}, p))(\vec{v}, q) := & - \int_I \int_{\Omega} \vec{u} \cdot \frac{\partial \vec{v}}{\partial t} dx dt + \int_I \int_{\Omega} \nu \nabla \vec{u} : \nabla \vec{v} dx dt \\ & + \int_I \int_{\Omega} \vec{v} \cdot \nabla p dx dt - \int_I \int_{\Omega} \vec{u} \cdot \nabla q dx dt \blacksquare \end{aligned}$$

Thm 3. *Let $\Omega \subset \mathbb{R}^n$ for $n \in \{2, 3\}$ either C^2 , or convex with a piecewise smooth bdr. Then with*

$$\mathcal{X} := L_2(I; L_2(\Omega)) \times \left(L_2(I; H^1(\Omega)/\mathbb{R}) \cap H_{0,\{T\}}^1(I; (H^1(\Omega)/\mathbb{R})') \right)'$$

$$\mathcal{Y} := L_2(I; (H_0^1(\Omega) \cap H^2(\Omega))^n) \cap H_{0,\{T\}}^1(I; L_2(\Omega)^n) \times L_2(I; H^1(\Omega)/\mathbb{R})$$

$B : \mathcal{X} \rightarrow \mathcal{Y}'$ is b.i.

Can be extended to NSE.

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