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RESEARCH ARTICLE

On some anisotropic, nonlocal, parabolic singular perturbations problems

Michel Chipot ^{a*} and Senoussi Guesmia^a

^aInstitute of Mathematics, University of Zurich, Winterthurerstrasse 190, CH-8057, Zurich, Switzerland. (Received 00 Month 200x; in final form 00 Month 200x)

This paper is devoted to the study of the anisotropic singular perturbations theory for some quasilinear parabolic problems. Describing the asymptotic behaviour of the solutions of non-local problems yields to show the existence of solution to some integro-differential problems. The closeness between the solutions of these two kinds of problems is established.

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1. Introduction

Let $\Omega = (0,1)^2$ be the unit square in \mathbb{R}^2 and T > 0 be a positive constant. If a is a positive, continuous function and $f \in L^2(0,T;L^2(\Omega))$, $u^0 \in L^2(\Omega)$ suppose that we are interested in solving the following integro-differential problem: find a function $u_0 = u_0(t, x_1, x_2)$ such that

$$\begin{cases} \partial_{t} u_{0}(\cdot, x_{1}, \cdot) - a\left(\int_{\Omega} u_{0}(\cdot, x) \, dx\right) \partial_{x_{2}}^{2} u_{0}(\cdot, x_{1}, \cdot) = f(\cdot, x_{1}, \cdot) \\ & \text{in } (0, T) \times (0, 1) \text{ a.e } x_{1} \in (0, 1) \,, \\ u_{0}(\cdot, x_{1}, \cdot) = 0 \text{ on } (0, T) \times \{0, 1\} \text{ a.e } x_{1} \in (0, 1) \,, \\ u_{0}(0, \cdot) = u^{0} \text{ in } \Omega. \end{cases}$$
(1)

Since the integral in the first equation of (1) is taken on the whole domain Ω it is clear that this equation cannot be considered as a parabolic equation parametrized by x_1 . Moreover, if starting from $w \in L^2(0,T; L^2(\Omega))$, one solves for u = Sw the problem parametrized by x_1

$$\begin{cases} \partial_t u \left(\cdot, x_1, \cdot \right) - a \left(\int_{\Omega} w \left(\cdot, x \right) dx \right) \partial_{x_2}^2 u \left(\cdot, x_1, \cdot \right) = f \left(\cdot, x_1, \cdot \right) \\ & \text{ in } (0, T) \times (0, 1) \text{ a.e } x_1 \in (0, 1) , \\ u \left(\cdot, x_1, \cdot \right) = 0 \text{ on } (0, T) \times \{0, 1\} \text{ a.e } x_1 \in (0, 1) , \\ u \left(0, x_1, \cdot \right) = u^0 \text{ in } (0, 1) \text{ a.e } x_1 \in (0, 1) , \end{cases}$$

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^{*}Corresponding author. Email: m.m.chipot@math.uzh.ch

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when it comes to finding a fixed point for S some compactness for this mapping is missing. One way to overcome this is to introduce for instance the following anisotropic singular perturbations problem

$$\begin{cases} \partial_t u_{\varepsilon} - a \left(\int_{\Omega} u_{\varepsilon} \left(\cdot, x \right) dx \right) \left(\varepsilon^2 \partial_{x_1}^2 u_{\varepsilon} + \partial_{x_2}^2 u_{\varepsilon} \right) = f \text{ in } (0, T) \times \Omega, \\ u_{\varepsilon} = 0 \text{ on } (0, T) \times \partial\Omega, \\ u_{\varepsilon} (0) = u^0 \text{ in } \Omega, \end{cases}$$

$$(2)$$

and to see what can be said on the behaviour of u_{ε} when $\varepsilon > 0$ goes to 0. This is what we would like to do in this note. For more details about the anisotropic singular perturbations problems from different point of view, the reader is referred to [1, 3–6, 8, 9] (see [10] for the basic theory of singular perturbations).

2. A more general setting

Now we suppose that Ω is a bounded open subset of \mathbb{R}^n and introduce a problem generalizing (2). We split the components of a point $x = (x_1, \ldots, x_n) \in \mathbb{R}^n$ into two parts

$$X_1 = (x_1, \dots, x_p)$$
 and $X_2 = (x_{p+1}, \dots, x_n)$,

so that $x = (X_1, X_2)$ (p and n are positive integers such that p < n). Similarly setting

$$\nabla_{X_1} u = (\partial_{x_1} u, \dots, \partial_{x_p} u)^T, \ \nabla_{X_2} u = (\partial_{x_{p+1}} u, \dots, \partial_{x_n} u)^T,$$

where T denotes the transpose operation and ∂_{x_i} the partial derivative in the direction x_i , we have if ∇ denotes the usual gradient operator

$$\nabla u = \begin{pmatrix} \nabla_{X_1} u \\ \nabla_{X_2} u \end{pmatrix}.$$
(3)

It will be convenient to simplify the notation to set

$$z = (t, x) \in \mathbb{R}^{n+1}, \ Q = (0, T) \times \Omega.$$
(4)

Let

$$A = A(z,s) = (a_{ij}(z,s)) \tag{5}$$

be a $n\times n$ matrix defined for $(z,s)\in Q\times \mathbb{R}$ and such that A is of Carathéodory type that is to say

$$z \mapsto a_{ij}(z,s)$$
 is measurable $\forall s \in \mathbb{R}, \quad \forall i, j = 1, \dots, n,$ (6)

$$s \mapsto a_{ij}(z,s)$$
 is continuous a.e. $z \in Q, \quad \forall i, j = 1, \dots, n.$ (7)

Moreover we assume that there exist two positive constants λ , Λ such that

$$\lambda |\xi|^2 \le A(z,s) \, \xi \cdot \xi \le \Lambda |\xi|^2 \quad \forall \xi \in \mathbb{R}^n, \text{ a.e. } z \in Q, \ \forall s \in \mathbb{R}.$$
(8)

In the inequality above we have denoted the scalar product by a dot and the usual Euclidean norm by $|\cdot|$. Note that (8) implies that the coefficients a_{ij} are uniformly bounded.

We split A into four blocks in such way that

$$A = \begin{pmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{pmatrix}, \tag{9}$$

where $A_{11},\,A_{22}$ are respectively $p\times p$ and $(n-p)\times (n-p)$ matrices. We then set for $\varepsilon>0$

$$A_{\varepsilon} = A_{\varepsilon}(z, s) = \begin{pmatrix} \varepsilon^2 A_{11} \ \varepsilon A_{12} \\ \varepsilon A_{21} \ A_{22} \end{pmatrix}.$$
 (10)

If $\xi = \begin{pmatrix} \bar{\xi}_1 \\ \bar{\xi}_2 \end{pmatrix} \in \mathbb{R}^n$ where $\bar{\xi}_1 = (\xi_1, \dots, \xi_p)^T$, $\bar{\xi}_2 = (\xi_{p+1}, \dots, \xi_n)^T$ by (8) one has for $\xi_{\varepsilon} = \begin{pmatrix} \varepsilon \bar{\xi}_1 \\ \bar{\xi}_2 \end{pmatrix}$

$$\lambda |\xi_{\varepsilon}|^{2} = \lambda \{ \varepsilon^{2} |\bar{\xi}_{1}|^{2} + |\bar{\xi}_{2}|^{2} \} \leq A \xi_{\varepsilon} \cdot \xi_{\varepsilon} = A_{\varepsilon} \xi \cdot \xi \quad \forall \xi \in \mathbb{R}^{n}, \text{ a.e. } z \in Q, \ \forall s \in \mathbb{R}, \ (11)$$
$$\lambda |\bar{\xi}_{2}|^{2} \leq A_{22} \bar{\xi}_{2} \cdot \bar{\xi}_{2} \quad \forall \bar{\xi}_{2} \in \mathbb{R}^{n-p}, \text{ a.e. } z \in Q, \ \forall s \in \mathbb{R}.$$
(12)

We would like to consider the family of parabolic, nonlinear, nonlocal problems

$$\begin{cases} \partial_t u_{\varepsilon} - \nabla \cdot \left(A_{\varepsilon}\left(\cdot, l\left(u_{\varepsilon}\right)\right) \nabla u_{\varepsilon}\right) = f & \text{in } Q, \\ u_{\varepsilon} = 0 & \text{on} \quad (0, T) \times \partial \Omega, \\ u_{\varepsilon}\left(0, \cdot\right) = u^0 & \text{in } \Omega, \end{cases}$$
(13)

where

$$l(u) = \int_{\Omega} h(x) u(t, x) dx$$
(14)

for some $h \in L^2(\Omega)$ i.e. l is a continuous linear form on $L^2(\Omega)$. We are in particular interested to show that u_{ε} - perhaps up to subsequence - possesses a limit when $\varepsilon \to 0$. Note that for A(z, s) = a(s) Id where Id is the identity matrix in \mathbb{R}^2 , $h \equiv 1$ one recognizes the problem (2) of the introduction and the limit of u_{ε} will provide us with a solution to problem (1).

We suppose here that

$$f \in L^2(Q), \ u^0 \in L^2(\Omega).$$

$$(15)$$

Then we have

Theorem 2.1: Under the assumptions above for any $\varepsilon > 0$ there exists a weak

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solution u_{ε} to problem (13) that is to say a function $u_{\varepsilon} = u_{\varepsilon}(t, x)$ such that

$$\begin{cases} u_{\varepsilon} \in L^{2}\left(0, T; H_{0}^{1}(\Omega)\right) \cap \mathcal{C}\left([0, T]; L^{2}(\Omega)\right), \ \partial_{t}u_{\varepsilon} \in L^{2}\left(0, T; H^{-1}(\Omega)\right), \\ \left\langle\partial_{t}u_{\varepsilon}, v\right\rangle + \int_{\Omega} A_{\varepsilon}\left(t, x, l\left(u_{\varepsilon}\right)\right) \nabla u_{\varepsilon} \cdot \nabla v \, dx = (f, v) \\ in \ \mathcal{D}'\left(0, T\right), \ \forall v \in H_{0}^{1}(\Omega), \\ u\left(0, \cdot\right) = u^{0}. \end{cases}$$

$$(16)$$

We refer the reader to [2], [7] for the different spaces introduced above. \mathcal{D}' denotes the space of distributions, $\langle \cdot, \cdot \rangle$ the duality bracket between $H^{-1}(\Omega)$ and $H_0^1(\Omega)$, (\cdot, \cdot) the usual scalar product in $L^2(\Omega)$.

Proof: The proof is a simple application of the Schauder fixed point theorem. Indeed for $w \in L^2(Q)$ one considers u = Sw the solution to

$$\begin{cases} u \in L^2\left(0, T; H_0^1(\Omega)\right) \cap \mathcal{C}\left([0, T]; L^2(\Omega)\right), \ \partial_t u \in L^2\left(0, T; H^{-1}(\Omega)\right), \\ \left\langle \partial_t u, v \right\rangle + \int_{\Omega} A_{\varepsilon}\left(t, x, l\left(w\right)\right) \nabla u \cdot \nabla v \, dx = (f, v) \quad \text{in } \mathcal{D}'\left(0, T\right), \ \forall v \in H_0^1(\Omega), \\ u\left(0, \cdot\right) = u^0. \end{cases}$$

Then it is easy to show (cf. [2]) that $w \to Sw$ has a fixed point u_{ε} . This completes the proof of the theorem.

3. Passage to the limit

We are now going to analyse the behaviour of u_{ε} when $\varepsilon \to 0$. For this purpose we introduce Ω_{X_1} the section of Ω above X_1 defined as

$$\Omega_{X_1} = \{ X_2 \mid (X_1, X_2) \in \Omega \}$$
(18)

(17)

and

$$Q_{X_1} = (0, T) \times \Omega_{X_1}.$$
 (19)

Then - if we let ε go to 0 very formally - we see that a reasonable candidate for the limit of u_{ε} when $\varepsilon \to 0$ is given by

$$u_0 = u_0(t, X_1, X_2) \tag{20}$$

where u_0 satisfies for a.e. $X_1 \in \Pi_1(\Omega) = \{X_1 \mid \exists X_2 \text{ such that } (X_1, X_2) \in \Omega \}$

$$\begin{aligned} \left\langle \partial_t u_0(\cdot, X_1, \cdot) - \nabla_{X_2} \cdot \left(A_{22} \left(\cdot, X_1, \cdot, l \left(u_0 \right) \right) \nabla_{X_2} u_0 \right) \left(\cdot, X_1, \cdot \right) \\ &= f(\cdot, X_1, \cdot) \text{ in } Q_{X_1}, \\ u_0(\cdot, X_1, \cdot) &= 0 \text{ on } (0, T) \times \partial \Omega_{X_1}, \\ u_0(0, X_1, \cdot) &= u^0 \left(X_1, \cdot \right) \text{ in } \Omega_{X_1}. \end{aligned}$$

$$\end{aligned}$$

$$(21)$$

Then we have

Theorem 3.1: Under the assumptions above there exists u_0 weak solution to (21) in the sense that a.e. $X_1 \in \Pi_1(\Omega)$

$$\begin{cases} u_{0}(\cdot, X_{1}, \cdot) \in L^{2}\left(0, T; H_{0}^{1}(\Omega_{X_{1}})\right) \cap \mathcal{C}\left([0, T]; L^{2}(\Omega_{X_{1}})\right), & u_{0} \in L^{2}\left(Q\right), \\ \partial_{t}u_{0}(\cdot, X_{1}, \cdot) \in L^{2}\left(0, T; H^{-1}(\Omega_{X_{1}})\right), \\ u_{0}\left(0, X_{1}, \cdot\right) = u^{0}\left(X_{1}, \cdot\right) \quad on \ \Omega_{X_{1}}, \\ \partial_{t}\left(u_{0}, v\right)_{\Omega_{X_{1}}} + \int_{\Omega_{X_{1}}} A_{22}\left(z, l\left(u_{0}\right)\right) \nabla_{X_{2}}u_{0} \cdot \nabla_{X_{2}}v \, dX_{2} = (f, v)_{\Omega_{X_{1}}} \\ \quad in \ \mathcal{D}'\left(0, T\right), \ \forall v \in H_{0}^{1}(\Omega_{X_{1}}). \end{cases}$$

$$(22)$$

 $(\cdot, \cdot)_{\Omega_{X_1}}$ denotes the usual $L^2(\Omega_{X_1})$ scalar product.

Remark 1: The asymptotic behaviour of the linear parabolic problems can be considered as a particular case. Indeed it is enough to choose A independent of s, i.e.

$$A(z,s) = A(z) \quad \forall z \in Q.$$

Proof: The solution u_0 is going to be obtained as a limit of u_{ε} when $\varepsilon \to 0$. That is to say we are going to show the following lemma which at the same time will complete the proof of Theorem 3.1.

Lemma 3.2: Under the assumptions above there exists $u_0 \in L^2(Q)$ and a "subsequence" of ε still labelled ε converging toward 0 such that

$$u_{\varepsilon} \rightharpoonup u_{0}, \qquad \varepsilon \nabla_{X_{1}} u_{\varepsilon} \rightharpoonup 0, \qquad \nabla_{X_{2}} u_{\varepsilon} \rightharpoonup \nabla_{X_{2}} u_{0} \quad in \ L^{2}(Q),$$
$$\partial_{t} u_{\varepsilon} \rightharpoonup \partial_{t} u_{0} \qquad in \ L^{2}\left(0, T; H^{-1}(\Omega)\right),$$

and u_0 is solution to (22).

(In the above the convergence in $L^2(\Omega)$ means for vectors the convergence of each components).

Proof: [Proof of the lemma 3.2]

If we take $v = u_{\varepsilon}$ in (16) -cf. [2], [7] for the technical details- we get for a.e. t

$$\frac{1}{2}\frac{d}{dt}\left|u_{\varepsilon}\right|^{2}_{2,\Omega} + \int_{\Omega} A_{\varepsilon}\left(t,\cdot,l\left(u_{\varepsilon}\right)\right) \nabla u_{\varepsilon} \cdot \nabla u_{\varepsilon} \, dx = \left(f,u_{\varepsilon}\right).$$

 $(|u_{\varepsilon}|_{2,\Omega})$ denotes the usual $L^2(\Omega)$ -norm). Using (11) and the Cauchy-Schwarz inequality we derive

$$\frac{1}{2}\frac{d}{dt}\left|u_{\varepsilon}\right|_{2,\Omega}^{2}+\lambda\int_{\Omega}\varepsilon^{2}|\nabla_{X_{1}}u_{\varepsilon}|^{2}+|\nabla_{X_{2}}u_{\varepsilon}|^{2}\,dx\leq(f,u_{\varepsilon})\leq|f\left(t,\cdot\right)|_{2,\Omega}|u_{\varepsilon}|_{2,\Omega}.$$

Since Ω is bounded -in particular in the directions X_2 - we have for some constant C > 0 independent of ε a Poincaré inequality of the type

$$|v|_{2,\Omega} \le C ||\nabla_{X_2} v||_{2,\Omega} \quad \forall v \in H_0^1(\Omega).$$

$$\tag{23}$$

 $(|\nabla_{X_2}v|$ denotes the Euclidean norm of $\nabla_{X_2}v$). Thus, by the Young inequality, we

deduce

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$$\frac{1}{2} \frac{d}{dt} |u_{\varepsilon}|^{2}_{2,\Omega} + \lambda \int_{\Omega} \varepsilon^{2} |\nabla_{X_{1}} u_{\varepsilon}|^{2} + |\nabla_{X_{2}} u_{\varepsilon}|^{2} dx \leq C |f|_{2,\Omega} ||\nabla_{X_{2}} u_{\varepsilon}||_{2,\Omega} \\
\leq \frac{C^{2}}{2\lambda} |f|^{2}_{2,\Omega} + \frac{\lambda}{2} ||\nabla_{X_{2}} u_{\varepsilon}||^{2}_{2,\Omega}$$

and in particular

$$\frac{1}{2}\frac{d}{dt}|u_{\varepsilon}|^{2}_{2,\Omega} + \frac{\lambda}{2}\int_{\Omega}\varepsilon^{2}|\nabla_{X_{1}}u_{\varepsilon}|^{2} + |\nabla_{X_{2}}u_{\varepsilon}|^{2}dx \le \frac{C^{2}}{2\lambda}|f|^{2}_{2,\Omega}.$$
(24)

Integrating in t between 0 and T we get

$$\frac{1}{2} |u_{\varepsilon}|_{2,\Omega}^{2}(T) + \frac{\lambda}{2} \int_{0}^{T} \int_{\Omega} \varepsilon^{2} |\nabla_{X_{1}} u_{\varepsilon}|^{2} + |\nabla_{X_{2}} u_{\varepsilon}|^{2} dt dx \leq \frac{1}{2} |u^{0}|_{2,\Omega}^{2} + \frac{C^{2}}{2\lambda} |f|_{2,Q}^{2}, \quad (25)$$

and thus

$$u_{\varepsilon}, \quad |\varepsilon \nabla_{X_1} u_{\varepsilon}|, \quad |\nabla_{X_2} u_{\varepsilon}| \quad \text{are bounded in } L^2(Q),$$
 (26)

(this of course independently of ε). From the equation in (16) we derive for $v \in H_0^1(\Omega)$,

$$\begin{aligned} |\langle \partial_t u_{\varepsilon}, v \rangle| &\leq |(f, v)| + \left| \int_{\Omega} A_{\varepsilon} \left(t, \cdot, l \left(u_{\varepsilon} \right) \right) \nabla u_{\varepsilon} \cdot \nabla v \, dx \right| \\ &\leq |f \left(t, \cdot \right)|_{2,\Omega} \left| v \right|_{2,\Omega} + C \left(\left| \left| \varepsilon \nabla_{X_1} u_{\varepsilon} \right| \left|_{2,\Omega} + \left| \left| \nabla_{X_2} u_{\varepsilon} \right| \left|_{2,\Omega} \right) \right| \left| \nabla_{X_2} v \right| \right|_{2,\Omega} \right) \\ &\leq C' \left(|f \left(t, \cdot \right)|_{2,\Omega} + \left| \left| \varepsilon \nabla_{X_1} u_{\varepsilon} \right| \left|_{2,\Omega} + \left| \left| \nabla_{X_2} u_{\varepsilon} \right| \right|_{2,\Omega} \right) \left| v \right|_{H^1(\Omega)}, \end{aligned}$$

where C, C' are constants independent of $\varepsilon \to 0$, $|v|_{H^1(\Omega)}$ the usual $H^1(\Omega)$ –norm. This implies

$$|\partial_t u_{\varepsilon}|_{L^2(0,T;H^{-1}(\Omega))} \le C' \left(|f|_{2,Q} + ||\varepsilon \nabla_{X_1} u_{\varepsilon}||_{2,Q} + ||\nabla_{X_2} u_{\varepsilon}||_{2,Q} \right).$$

It follows from (26) that

$$\partial_t u_{\varepsilon}$$
 is bounded in $L^2(0,T;H^{-1}(\Omega))$. (27)

Up to a subsequence we deduce from (26), (27) that there exist $u_0, u_1, u_2 \in L^2(Q)$ and $u_3 \in L^2(0, T; H^{-1}(\Omega))$ - i.e. with components in $L^2(Q)$ for the vectors - such that

$$u_{\varepsilon} \rightharpoonup u_0, \quad \varepsilon \nabla_{X_1} u_{\varepsilon} \rightharpoonup u_1, \quad \nabla_{X_2} u_{\varepsilon} \rightharpoonup u_2, \quad \text{in } L^2(Q),$$

 $\partial_t u_{\varepsilon} \rightharpoonup u_3 \quad \text{in } L^2(0,T; H^{-1}(\Omega)).$

By the continuity of the derivation in $\mathcal{D}'(Q)$ we derive that

$$u_{\varepsilon} \rightharpoonup u_0, \qquad \varepsilon \nabla_{X_1} u_{\varepsilon} \rightharpoonup 0, \qquad \nabla_{X_2} u_{\varepsilon} \rightharpoonup \nabla_{X_2} u_0 \quad \text{in } L^2(Q),$$
 (28)

$$\partial_t u_{\varepsilon} \rightharpoonup \partial_t u_0 \quad \text{in } L^2(0,T; H^{-1}(\Omega)).$$
 (29)

A delicate issue is now to show that

$$u_0(t, X_1, \cdot) \in H^1_0(\Omega_{X_1})$$
 a.e. $(t, X_1) \in (0, T) \times \Pi_1(\Omega)$.

For that we denote by $\mathcal B$ an open ball of $\mathbb R^{n-p}$ such that

$$\Omega_{X_1} \subset \mathcal{B} \quad \forall X_1 \in \Pi_1(\Omega) \,. \tag{30}$$

We suppose that $u_{\varepsilon}(t, \cdot)$ is extended by 0 outside Ω . Then from (26) we derive that

$$\int_0^T \int_{\Pi_1(\Omega) \times \mathcal{B}} \left| \nabla_{X_2} u_{\varepsilon} \right|^2 dt dx \le C,\tag{31}$$

where C is a constant independent of ε . This can also be written as

$$|u_{\varepsilon}|_{L^2((0,T)\times\Pi_1(\Omega);H^1_0(\mathcal{B}))} \le C.$$

Thus there exists a function $\bar{u}_0 \in L^2((0,T) \times \Pi_1(\Omega); H_0^1(\mathcal{B}))$, a subsequence of the subsequence above and still labelled ε such that

$$u_{\varepsilon} \rightharpoonup \bar{u}_{0} \quad \text{in } L^{2}\left((0,T) \times \Pi_{1}\left(\Omega\right); H^{1}_{0}\left(\mathcal{B}\right)\right)$$

and in particular

$$u_{\varepsilon} \to \bar{u}_0 \quad \text{in } \mathcal{D}'\left((0,T) \times \Pi_1\left(\Omega\right) \times \mathcal{B}\right)$$

If u_0 is also extended by 0 outside Ω one has of course also

$$u_{\varepsilon} \to u_0 \quad \text{in } \mathcal{D}' \left((0, T) \times \Pi_1 \left(\Omega \right) \times \mathcal{B} \right)$$

and thus

$$u_0 = \bar{u}_0 \in L^2\left((0,T) \times \Pi_1\left(\Omega\right); H^1_0\left(\mathcal{B}\right)\right).$$

It follows that

$$u_0(t, X_1, \cdot) \in H_0^1(\mathcal{B})$$
 a.e. $(t, X_1) \in (0, T) \times \Pi_1(\Omega)$

i.e.

$$u_0(t, X_1, \cdot) \in H^1_0(\Omega_{X_1})$$
 a.e. $(t, X_1) \in (0, T) \times \Pi_1(\Omega)$. (32)

We next show that

$$l(u_{\varepsilon}) \to l(u_0)$$
.

This will follow from the lemma:

Lemma 3.3: For any $v \in H_0^1(\Omega)$, the functions

$$t \mapsto \int_{\Omega} u_{\varepsilon} v dx, \quad t \mapsto \int_{\Omega} u_0 v dx$$
 (33)

belong to $H^{1}\left(0,T\right)$ and for the subsequence above we have

(i)
$$\int_{\Omega} u_{\varepsilon} v dx \to \int_{\Omega} u_0 v dx \quad in \ L^2(0,T) , \ \mathcal{C}(0,T) ,$$

(ii)
$$\int_{\Omega} u_{\varepsilon} v dx \rightharpoonup \int_{\Omega} u_0 v dx \quad in \ H^1(0,T) ,$$

(iii)
$$l(u_{\varepsilon}) \to l(u_0) \quad in \ L^2(0,T) .$$

 $(\mathcal{C}(0,T) \text{ denotes the space of continuous functions on } [0,T] \text{ for the uniform norm}).$

Proof: [Proof of the lemma 3.3]

From (16) we derive for $v \in H_0^1(\Omega)$

$$\frac{d}{dt}\left(u_{\varepsilon}\left(t,\cdot\right),v\right) = \left(f\left(t,\cdot\right),v\right) - \int_{\Omega} A_{\varepsilon}\left(t,x,l\left(u_{\varepsilon}\right)\right) \nabla u_{\varepsilon}\left(t,x\right) \cdot \nabla v\left(x\right) \, dx.$$

Thus we have (see (8)) for some constant C independent of ε

$$\left|\frac{d}{dt}\left(u_{\varepsilon},v\right)\right| \leq \left|f\right|_{2,\Omega}\left|v\right|_{2,\Omega} + C\left(\varepsilon\right)\left|\nabla_{X_{1}}u_{\varepsilon}\right|\left|_{2,\Omega} + \left|\left|\nabla_{X_{2}}u_{\varepsilon}\right|\right|_{2,\Omega}\right)\left|\left|\nabla v\right|\right|_{2,\Omega}, \quad (34)$$

whence

$$\frac{d}{dt}\left(u_{\varepsilon},v\right)\in L^{2}\left(0,T\right)$$

Similarly one has

$$|(u_{\varepsilon}, v)| \le |u_{\varepsilon}|_{2,\Omega} |v|_{2,\Omega}$$

and by (26), (34) we conclude that

$$(u_{\varepsilon}, v) \in H^1(0, T)$$

and for some constant D independent of ε it holds

$$|(u_{\varepsilon}, v)|_{H^1(0,T)} \le D.$$

Thus - up to a subsequence of the above subsequence - we have

$$(u_{\varepsilon}, v) \to L_v \text{ in } L^2(0, T), \ \mathcal{C}(0, T),$$

$$(35)$$

$$(u_{\varepsilon}, v) \rightharpoonup L_v \text{ in } H^1(0, T).$$
 (36)

Let us choose $\varphi \in \mathcal{D}(0,T)$ then one has

$$\int_0^T \left(u_{\varepsilon}, v\right) \varphi dt = \int_0^T \int_{\Omega} u_{\varepsilon} v \varphi dt dx \to \int_0^T \int_{\Omega} u_0 v \varphi dt dx = \int_0^T \left(u_0, v\right) \varphi dt,$$

by (28). On the other hand by (35) we have

$$\int_0^T \left(u_\varepsilon, v \right) \varphi dt \to \int_0^T L_v \varphi dt.$$

It follows that

$$L_v = \int_{\Omega} u_0 v dx$$

and by uniqueness of the limit, the convergences in (35), (36) are not up to a subsequence. This completes the proof of (i), (ii) and (iii) for $h \in H_0^1(\Omega)$ (recall (14)). To obtain (iii) for $h \in L^2(\Omega)$ it is enough to rely in the density of $H_0^1(\Omega)$ in $L^2(\Omega)$ taking into account (26).

End of the proof of Lemma 3.2 and Theorem 3.1 For $\varphi \in \mathcal{D}(0,T), v \in H_0^1(\Omega)$ we have by (16)

$$-\int_{0}^{T} (u_{\varepsilon}, v) \varphi' dt + \int_{0}^{T} \int_{\Omega} \varepsilon A_{11} (z, l(u_{\varepsilon})) \varepsilon \nabla_{X_{1}} u_{\varepsilon} \cdot \nabla_{X_{1}} v \varphi dt dx + \int_{0}^{T} \int_{\Omega} A_{12} (z, l(u_{\varepsilon})) \varepsilon \nabla_{X_{2}} u_{\varepsilon} \cdot \nabla_{X_{1}} v \varphi dt dx + \int_{0}^{T} \int_{\Omega} A_{21} (z, l(u_{\varepsilon})) \varepsilon \nabla_{X_{1}} u_{\varepsilon} \cdot \nabla_{X_{2}} v \varphi dt dx + \int_{0}^{T} \int_{\Omega} A_{22} (z, l(u_{\varepsilon})) \nabla_{X_{2}} u_{\varepsilon} \cdot \nabla_{X_{2}} v \varphi dt dx = \int_{0}^{T} \int_{\Omega} f v \varphi dt dx.$$
(37)

Noting (iii) of Lemma 3.3 we have -up to a subsequence-

$$A(z, l(u_{\varepsilon})) \to A(z, l(u_0))$$
 a.e. $z \in Q$, (38)

i.e. we have this convergence for all coefficients of A. It follows that for $v \in H_0^1(\Omega)$

$$a_{ij}(\cdot, l(u_{\varepsilon})) \partial_{x_k} v \to a_{ij}(\cdot, l(u_0)) \partial_{x_k} v \text{ in } L^2(Q), \quad \forall i, j, k = 1, \dots, n.$$

Then by (28), (29) one can easily pass to the limit in (37) to get

$$-\int_{0}^{T} (u_{0}, v) \varphi' dt + \int_{0}^{T} \int_{\Omega} A_{22} (z, l(u_{0})) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} v \varphi dt dx = \int_{0}^{T} \int_{\Omega} f v \varphi dt dx.$$
(39)

Using an argument as in [3] or [5] one can easily conclude that

$$\int_{0}^{T} (u_{0}, v)_{\Omega_{X_{1}}} \varphi' dt + \int_{0}^{T} \int_{\Omega_{X_{1}}} A_{22} (t, X_{1}, X_{2}, l(u_{0})) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} v \varphi dt dx$$
$$= \int_{0}^{T} \int_{\Omega_{X_{1}}} f (t, X_{1}, X_{2}) v \varphi dt dx$$
$$\forall \varphi \in \mathcal{D} (0, T), \quad \forall v \in H_{0}^{1}(\Omega_{X_{1}}), \text{ a.e. } X_{1} \in \Pi_{1}(\Omega),$$
(40)

which is the last equation in (22). (We denoted by $(\cdot, \cdot)_{\Omega_{X_1}}$ the scalar product in $L^2(\Omega_{X_1})$). Clearly (40) implies that for a.e. $X_1 \in \Pi_1(\Omega)$

$$\partial_t u_0 = \nabla_{X_2} \left(A_{22} \left(\cdot, l \left(u_0 \right) \right) \nabla_{X_2} u_0 \right) + f \quad \text{in } \mathcal{D}' \left((0, T) \times \Omega_{X_1} \right)$$
(41)

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and since $u_0 \in L^2\left((0,T); H_0^1\left(\Omega_{X_1}\right)\right)$ we derive

$$\partial_t u_0 \in L^2(0,T; H^{-1}(\Omega_{X_1}))$$
 for a.e. $X_1 \in \Pi_1(\Omega)$

Since for a.e. $X_1 \in \Pi_1(\Omega)$

$$u_0 \in L^2(0,T; L^2(\Omega_{X_1}))$$
 (42)

by the usual embedding theorem the two first relations of (22) follow.

For $v \in H_0^1(\Omega)$ we have when $\varepsilon \to 0$

$$\left(u^{0},v\right)=\left(u_{\varepsilon}\left(0,\cdot\right),v\right)=\int_{\Omega}u_{\varepsilon}\left(0,\cdot\right)vdx\rightarrow\int_{\Omega}u_{0}\left(0,\cdot\right)vdx$$

by (i) of Lemma 3.3. Thus

$$(u^0, v) = (u_0(0, \cdot), v) \quad \forall v \in H_0^1(\Omega).$$
 (43)

By density of $H_0^1(\Omega)$ in $L^2(\Omega)$, (43) holds for every $v \in L^2(\Omega)$ and we get

$$u_0\left(0,\cdot\right) = u^0.$$

This completes the proof of Lemma 3.2.

Of course the proof of Theorem 3.1 is also completed.

4. Additional results

We now give additional properties of the solutions to (22). We have indeed

Theorem 4.1: Let u_0 be solution to (22), then

$$u_{0} \in V = \left\{ v \in L^{2}(Q) \mid |\nabla_{X_{2}}v| \in L^{2}(Q), \partial_{t}v \in L^{2}(0,T;H^{-1}(\Omega)) \right\}, \quad (44)$$

$$u_0 \in L^{\infty}\left(0, T; L^2\left(\Omega\right)\right). \tag{45}$$

Proof: From (22) we have for a.e. $X_1 \in \Pi_1(\Omega), v \in H^1_0(\Omega_{\Omega_{X_1}})$

$$\langle \partial_{t} u_{0}, v \rangle_{\Omega_{X_{1}}} + \int_{\Omega_{X_{1}}} A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} v \, dX_{2} = \int_{\Omega_{X_{1}}} f v dX_{2}$$

 $(\langle \cdot, \cdot \rangle_{\Omega_{X_1}}$ denotes the duality bracket between $H^{-1}(\Omega_{X_1})$ and $H^1_0(\Omega_{X_1})$). Choosing $v = u_0$ we deduce

$$\frac{1}{2}\frac{d}{dt}|u_0|^2_{2,\Omega_{X_1}} + \int_{\Omega_{X_1}} A_{22}(z,l(u_0))\nabla_{X_2}u_0 \cdot \nabla_{X_2}u_0 dX_2 = \int_{\Omega_{X_1}} fu_0 dX_2.$$

Integrating on $(0, t_0) \times \Pi_1(\Omega)$ we get

$$\frac{1}{2} |u_0|_{2,\Omega}^2 (t_0) + \int_0^{t_0} \int_{\Omega} A_{22} (z, l(u_0)) \nabla_{X_2} u_0 \cdot \nabla_{X_2} u_0 dx$$
$$= \frac{1}{2} |u^0|_{2,\Omega}^2 + \int_0^{t_0} \int_{\Omega} f u_0 \, dt dx.$$
(46)

Choosing $t_0 = T$ and using (12) we obtain

$$\lambda \| \nabla_{X_2} u_0 \|_{2,Q}^2 \le \frac{1}{2} \| u^0 \|_{2,\Omega}^2 + \| f \|_{2,Q} \| u_0 \|_{2,Q}.$$

Applying the Poincaré inequality on each section Ω_{X_1} we obtain for a constant C

$$|u_0|_{2,Q}^2 \le C ||\nabla_{X_2} u_0||_{2,Q}^2$$

and thus by Young's inequality

$$\lambda ||\nabla_{X_2} u_0||_{2,Q}^2 \le \frac{1}{2} |u^0|_{2,\Omega}^2 + \frac{C^2}{2\lambda} |f|_{2,Q}^2 + \frac{\lambda}{2} ||\nabla_{X_2} u_0||_{2,Q}^2$$

From this it follows that

$$||
abla_{X_2} u_0||^2_{2,Q} \le rac{1}{\lambda} ||u^0|^2_{2,\Omega} + rac{C^2}{\lambda^2} |f|^2_{2,Q}$$

On the other hand integrating the equation in (22) on $\Pi_1(\Omega)$ we have for $v \in H_0^1(\Omega)$

$$\begin{aligned} \langle \partial_t u_0, v \rangle &= \frac{d}{dt} \left(u_0, v \right) = \int_{\Pi_1(\Omega)} \partial_t \left(u_0, v \right)_{\mathcal{B}} dX_1 = \int_{\Pi_1(\Omega)} \partial_t \left(u_0, v \right)_{\Omega_{X_1}} dX_1 \\ &= \int_{\Omega} f v \, dx - \int_{\Omega} A_{22} \left(\cdot, l \left(u_0 \right) \right) \nabla_{X_2} u_0 \cdot \nabla_{X_2} v \, dx. \end{aligned}$$

where \mathcal{B} is defined by (30) and u_0 , v are supposed to be extended by 0 outside Ω . Then it follows that

$$\begin{aligned} |\langle \partial_t u_0, v \rangle| &\leq |(f, v)| + \left| \int_{\Omega} A_{22} \left(\cdot, l \left(u_0 \right) \right) \nabla u_0 \cdot \nabla v \, dx \right| \\ &\leq C \left(|f \left(t, X_1, \cdot \right)|_{2,\Omega} + ||\nabla_{X_2} u_0||_{2,\Omega} \right) |v|_{H^1(\Omega)} \,, \end{aligned}$$

where C is a positive constant. Then after integrating on (0, T) we get

$$|\partial_t u_0|_{L^2(0,T;H^{-1}(\Omega))} \le C\left(|f|_{2,Q} + ||\nabla_{X_2} u_0||_{2,Q}\right)$$

Thus (44) follows and (45) is then a consequence of (46).

In the case where the solution of (22) is unique we can yet prove

Theorem 4.2: Suppose that (22) has a unique solution then we have

$$u_{\varepsilon} \longrightarrow u_0 \quad in \ V \tag{47}$$

i.e.

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$$u_{\varepsilon} \longrightarrow u_0, \quad \nabla_{X_2} u_{\varepsilon} \longrightarrow \nabla_{X_2} u_0 \quad in \ L^2(Q), \quad \partial_t u_{\varepsilon} \longrightarrow \partial_t u_0 \quad in \ L^2(0,T; H^{-1}(\Omega))$$

and

$$\varepsilon \nabla_{X_1} u_{\varepsilon} \longrightarrow 0 \quad in \ L^2(Q).$$
 (48)

(The convergences hold for the whole sequence).

Proof: We introduce

$$I_{\varepsilon} = \int_{Q} A_{\varepsilon} \left(z, l\left(u_{\varepsilon} \right) \right) \begin{pmatrix} \nabla_{X_{1}} u_{\varepsilon} \\ \nabla_{X_{2}} \left(u_{\varepsilon} - u_{0} \right) \end{pmatrix} \cdot \begin{pmatrix} \nabla_{X_{1}} u_{\varepsilon} \\ \nabla_{X_{2}} \left(u_{\varepsilon} - u_{0} \right) \end{pmatrix} dt dx.$$

In what follows we just denote by A_{ε} the matrix $A_{\varepsilon}(z, l(u_{\varepsilon}))$ and drop the measures of integration to get

$$\begin{split} I_{\varepsilon} &= \int_{Q} A_{\varepsilon} \nabla u_{\varepsilon} \cdot \nabla u_{\varepsilon} - \int_{Q} A_{\varepsilon} \nabla u_{\varepsilon} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} - \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \nabla u_{\varepsilon} \\ &+ \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \\ &= \int_{Q} f u_{\varepsilon} - \int_{0}^{T} \langle \partial_{t} u_{\varepsilon}, u_{\varepsilon} \rangle - \int_{Q} A_{\varepsilon} \nabla u_{\varepsilon} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} - \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \nabla u_{\varepsilon} \\ &+ \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \\ &:= \int_{Q} f u_{\varepsilon} - \int_{0}^{T} \langle \partial_{t} u_{\varepsilon}, u_{\varepsilon} \rangle dt + J_{\varepsilon}. \end{split}$$
(49)

(We used (16) with $v = u_{\varepsilon}$). Remark now that

$$\int_{Q} A_{\varepsilon} \nabla u_{\varepsilon} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} = \int_{Q} A_{21} \left(z, l \left(u_{\varepsilon} \right) \right) \varepsilon \nabla_{X_{1}} u_{\varepsilon} \cdot \nabla_{X_{2}} u_{0} + \int_{Q} A_{22} \left(z, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{2}} u_{\varepsilon} \cdot \nabla_{X_{2}} u_{0}.$$

Passing to the limit as we did in (37) we obtain that

$$\lim_{\varepsilon \to 0} \int_{Q} A_{\varepsilon} \nabla u_{\varepsilon} \cdot \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} = \int_{Q} A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} u_{0}.$$

Note that since the possible limit is unique this is not up to a subsequence. We obtain similarly

$$\lim_{\varepsilon \to 0} \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \nabla u_{\varepsilon} = \int_{Q} A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} u_{0},$$
$$\lim_{\varepsilon \to 0} \int_{Q} A_{\varepsilon} \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} \cdot \nabla \begin{pmatrix} 0 \\ \nabla_{X_{2}} u_{0} \end{pmatrix} = \int_{Q} A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} u_{0},$$

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and thus combining these three passages to the limit

$$\lim_{\varepsilon \to 0} J_{\varepsilon} = -\int_{Q} A_{22}\left(z, l\left(u_{0}\right)\right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} u_{0}.$$
(50)

We now consider the other terms in I_{ε} -see (49)- in order to pass to the limit. Let $\theta_k \in \mathcal{D}(\Omega)$ be a sequence of smooth functions independent of t that we will choose later on. We have

$$\begin{split} &\int_0^T \langle \partial_t u_{\varepsilon}, u_{\varepsilon} \rangle dt \\ &= \int_0^T \langle \partial_t \left(u_{\varepsilon} - \theta_k \right), u_{\varepsilon} - \theta_k \rangle dt + \int_0^T \langle \partial_t \theta_k, u_{\varepsilon} - \theta_k \rangle dt + \int_0^T \langle \partial_t u_{\varepsilon}, \theta_k \rangle dt \\ &= \frac{1}{2} \left| u_{\varepsilon} - \theta_k \right|_{2,\Omega}^2 (T) - \frac{1}{2} \left| u^0 - \theta_k \right|_{2,\Omega}^2 + \int_0^T \langle \partial_t u_{\varepsilon}, \theta_k \rangle dt. \end{split}$$

Going back to (49) we have

$$R_{\varepsilon} := I_{\varepsilon} + \frac{1}{2} \left| u_{\varepsilon} - \theta_k \right|_{2,\Omega}^2 (T) = \frac{1}{2} \left| u^0 - \theta_k \right|_{2,\Omega}^2 + \int_Q f u_{\varepsilon} - \int_0^T \langle \partial_t u_{\varepsilon}, \theta_k \rangle dt + J_{\varepsilon}.$$
(51)

Using (22), (28), (29), (50) we get

$$\begin{split} &\lim_{\varepsilon \to 0} R_{\varepsilon} \\ &= \frac{1}{2} \left| u^{0} - \theta_{k} \right|_{2,\Omega}^{2} + \int_{Q} f u_{0} - \int_{Q} A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} u_{0} - \int_{0}^{T} \langle \partial_{t} u_{0}, \theta_{k} \rangle dt \\ &= \frac{1}{2} \left| u^{0} - \theta_{k} \right|_{2,\Omega}^{2} + \int_{0}^{T} \int_{\Pi_{1}(\Omega)} \langle \partial_{t} u_{0}, u_{0} \rangle_{\Omega_{X_{1}}} \left(t, X_{1} \right) dt - \int_{0}^{T} \langle \partial_{t} u_{0}, \theta_{k} \rangle dt. \end{split}$$

Note that since we have

$$\int_0^T \langle \partial_t u_0, \theta_k \rangle dt = \int_0^T \int_{\Pi_1(\Omega)} \langle \partial_t u_0, \theta_k \rangle_{\Omega_{X_1}}(t, X_1) \, dt dX_1$$

the limit above becomes

$$\begin{split} \lim_{\varepsilon \to 0} R_{\varepsilon} &= \frac{1}{2} \left| u^{0} - \theta_{k} \right|_{2,\Omega}^{2} + \int_{\Pi_{1}(\Omega)} \int_{0}^{T} \langle \partial_{t} \left(u_{0} - \theta_{k} \right), \left(u_{0} - \theta_{k} \right) \rangle_{\Omega_{X_{1}}} \left(t, X_{1} \right) dt dX_{1} \\ &= \frac{1}{2} \left| u^{0} - \theta_{k} \right|_{2,\Omega}^{2} + \frac{1}{2} \int_{\Pi_{1}(\Omega)} \left| u_{0} - \theta_{k} \right|_{2,\Omega_{X_{1}}}^{2} \left(T, X_{1} \right) - \left| u^{0} - \theta_{k} \right|_{2,\Omega_{X_{1}}}^{2} \left(X_{1} \right) dt dX_{1} \\ &= \frac{1}{2} \left| u_{0} - \theta_{k} \right|_{2,\Omega}^{2} \left(T \right). \end{split}$$

$$(52)$$

Thus from (51) we deduce

$$\limsup_{\varepsilon \to 0} I_{\varepsilon} \leq \lim_{\varepsilon \to 0} R_{\varepsilon} = \frac{1}{2} |u_0 - \theta_k|^2_{2,\Omega}(T).$$

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Choosing θ_k such that

$$\theta_k \to u_0(T) \quad \text{in } L^2(\Omega),$$
(53)

we deduce that

$$0 \le \limsup_{\varepsilon \to 0} I_{\varepsilon} \le 0$$

and thus

$$I_{\varepsilon} \to 0 \quad \text{when } \varepsilon \to 0$$

Since, by the coerciveness assumption, we have

$$\lambda \int_0^T \int_\Omega |\varepsilon \nabla_{X_1} u_\varepsilon|^2 + |\nabla_{X_2} (u_\varepsilon - u_0)|^2 \le I_\varepsilon,$$

it follows that

$$\varepsilon \nabla_{X_1} u_{\varepsilon} \longrightarrow 0, \qquad \nabla_{X_2} u_{\varepsilon} \longrightarrow \nabla_{X_2} u_0 \quad \text{in } L^2(Q),$$

and thus (by Poincaré's inequality)

$$u_{\varepsilon} \longrightarrow u_0$$
 in $L^2(Q)$.

It remains to show that

$$\partial_t u_{\varepsilon} \longrightarrow \partial_t u_0 \quad \text{in } L^2\left(0, T; H^{-1}\left(\Omega\right)\right).$$

Testing the equation in (22) with $v \in H_0^1(\Omega)$ and integrating on $\Pi_1(\Omega)$, we get

$$\int_{\Pi_{1}(\Omega)} \partial_{t} (u_{0}, v)_{\Omega_{X_{1}}} dX_{1} + \int_{\Omega} A_{22} (z, l(u_{0})) \nabla_{X_{2}} u_{0} \cdot \nabla_{X_{2}} v \, dx = \int_{\Omega} f v \, dx,$$

whence

$$\frac{d}{dt} (u_0, v)_{\Omega} + \int_{\Omega} A_{22} (z, l(u_0)) \nabla_{X_2} u_0 \cdot \nabla_{X_2} v \, dx = \int_{\Omega} f v \, dx.$$

Then subtracting the above equality from (16) leads to

$$\begin{aligned} \frac{d}{dt} \left(u_{\varepsilon} - u_{0}, v \right) &+ \int_{\Omega} \varepsilon^{2} A_{11} \left(z, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{1}} u_{\varepsilon} \cdot \nabla_{X_{1}} v \, dx \\ &+ \int_{\Omega} \varepsilon A_{12} \left(z, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{2}} u_{\varepsilon} \cdot \nabla_{X_{1}} v \, dx \\ &+ \int_{\Omega} \varepsilon A_{21} \left(z, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{1}} u_{\varepsilon} \cdot \nabla_{X_{2}} v \, dx \\ &+ \int_{\Omega} \left(A_{22} \left(z, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{2}} u_{\varepsilon} - A_{22} \left(z, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \right) \cdot \nabla_{X_{2}} v \, dx \\ &= 0. \end{aligned}$$

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Using the Cauchy-Schwarz inequality and by (44), it follows that

$$\begin{aligned} \left\langle \partial_{t} \left(u_{\varepsilon} - u_{0} \right), v \right\rangle &\leq \varepsilon^{2} C \left| \left| \nabla_{X_{1}} u_{\varepsilon} \right| \right|_{2,\Omega} \left| \left| \nabla_{X_{1}} v \right| \right|_{2,\Omega} \\ &+ \varepsilon C \left| \left| \nabla_{X_{2}} u_{\varepsilon} \right| \right|_{2,\Omega} \left| \left| \nabla_{X_{1}} v \right| \right|_{2,\Omega} + \varepsilon C \left| \left| \nabla_{X_{1}} u_{\varepsilon} \right| \right|_{2,\Omega} \right| \left| \nabla_{X_{2}} v \right| \right|_{2,\Omega} \\ &+ \left| \left| A_{22} \left(\cdot, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{2}} u_{\varepsilon} - A_{22} \left(\cdot, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \right| \left| 2,\Omega \right| \left| \nabla_{X_{1}} v \right| \right|_{2,\Omega} \\ &\leq \left(C \left(\varepsilon^{2} \right) \left| \left| \nabla_{X_{1}} u_{\varepsilon} \right| \right|_{2,\Omega} + \varepsilon \left| \left| \nabla_{X_{2}} u_{\varepsilon} \right| \right|_{2,\Omega} + \varepsilon \left| \left| \nabla_{X_{1}} u_{\varepsilon} \right| \right|_{2,\Omega} \right) \\ &+ \left| \left| A_{22} \left(\cdot, l \left(u_{\varepsilon} \right) \right) \nabla_{X_{2}} u_{\varepsilon} - A_{22} \left(\cdot, l \left(u_{0} \right) \right) \nabla_{X_{2}} u_{0} \right| \left| 2,\Omega \right) \left| v \right|_{H^{1}(\Omega)}, \end{aligned}$$

for any $v \in H_0^1(\Omega)$. This implies

$$\begin{aligned} |\partial_t \left(u_{\varepsilon} - u_0 \right)|_{H^{-1}(\Omega)} &\leq C \left(\varepsilon^2 ||\nabla_{X_1} u_{\varepsilon}||_{2,\Omega} + \varepsilon ||\nabla_{X_2} u_{\varepsilon}||_{2,\Omega} + \varepsilon ||\nabla_{X_1} u_{\varepsilon}||_{2,\Omega} \right) \\ &+ ||A_{22} \left(\cdot, l \left(u_{\varepsilon} \right) \right) \nabla_{X_2} u_{\varepsilon} - A_{22} \left(\cdot, l \left(u_0 \right) \right) \nabla_{X_2} u_0 ||_{2,\Omega}. \end{aligned}$$

Integrating over (0, T), letting $\varepsilon \to 0$ and taking into account *(iii)* of Lemma 3.3, we deduce

$$|\partial_t (u_{\varepsilon} - u_0)|_{L^2(0,T;H^{-1}(\Omega))} \to 0.$$

This completes the proof of the theorem.

Remark 1: We also have

$$u_{\varepsilon}(t,\cdot) \longrightarrow u_0(t,\cdot) \quad \text{in } L^2(\Omega) \quad \forall t \in (0,T).$$
 (54)

Indeed, if we replace T by $t_0 \in (0, T)$ in the proof of Theorem 4.2 we get from (51), (52),

$$\limsup_{\varepsilon \to 0} \left\| \theta_k - u_{\varepsilon} \left(t_0, \cdot \right) \right\|_{2,\Omega} \le \left\| \theta_k - u_0 \left(t_0, \cdot \right) \right\|_{2,\Omega},$$

 ${\rm thus}$

$$\limsup_{\varepsilon \to 0} |u_{\varepsilon}(t_0, \cdot) - u_0(t_0, \cdot)|_{2,\Omega} \le |\theta_k - u_0(t_0, \cdot)|_{2,\Omega} + \limsup_{\varepsilon \to 0} |\theta_k - u_{\varepsilon}(t_0, \cdot)|_{2,\Omega}$$
$$\le 2 |\theta_k - u_0(t_0, \cdot)|_{2,\Omega}.$$
(55)

Then passing to the limit in the inequality above when $k \to \infty$ and taking into account (53) for $T = t_0$, we obtain

$$\limsup_{\varepsilon \to 0} \left| u_{\varepsilon} \left(t_0, \cdot \right) - u_0 \left(t_0, \cdot \right) \right|_{2,\Omega} \le 0,$$

whence

$$\lim_{\varepsilon \to 0} |u_{\varepsilon}(t_0, \cdot) - u_0(t_0, \cdot)|_{2,\Omega} = 0.$$

It follows that

$$u_{\varepsilon}(t,\cdot) \longrightarrow u_0(t,\cdot) \quad \text{in } L^2(\Omega) \quad \forall t \in (0,T) .$$

REFERENCES

Remark 2: In the case when u_0 is not unique one can show that the set of solutions to (22) lies in a neighborhood of

$$\{u_{\varepsilon} \mid \varepsilon < \varepsilon_0\}$$

for the topology of V, that is to say from any subsequence of u_{ε} there exists another subsequence such that (47), (48) hold. In addition if there exists a subsequence of u_{ε} , still labeled by u_{ε} , and a function u_0 such that

$$u_{\varepsilon} \longrightarrow u_0 \quad \text{in } \mathcal{D}'(Q),$$

then u_0 is a solution of (22) and (47), (48) hold.

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References

- B. Brighi, S. Guesmia, Asymptotic behavior of solutions of hyperbolic problems on a cylindrical domain, Discrete Contin. Dyn. Syst., suppl. (2007), pp. 160 - 169.
- [2] M. Chipot, Element of nonlinear analysis, Birkhäuser, 2000.
- [3] M. Chipot, Elliptic Equations: An Introductory Course, Birkhäuser, 2009.
- M. Chipot, On some anisotropic singular perturbation problems, Asymptotic Anal. 55 (3-4) (2007), pp. 125-144.
- [5] M. Chipot, S. Guesmia., On the asymptotic behavior of elliptic, anisotropic singular perturbations problems, Commun. Pure Appl. Anal. 8 (1) (2009), pp. 179-193.
- [6] M. Chipot, S. Guesmia, Correctors for some asymptotic problems, (submitted).
- [7] R. Dautray, J. L. Lions, Analyse mathématique et calcul numérique, volume 3, Masson, 1987.
- [8] S. Guesmia, Etude du comportement asymptotique de certaines équations aux dérivées partielles dans des domaines cylindriques. Thèse Université de Haute Alsace, December 2006.
- [9] S. Guesmia, On the asymptotic behavior of elliptic boundary value problems with some small coefficients, Electro. J. Differ. Equ. 59 (2008), pp. 1-13.
- [10] J. L. Lions, Perturbations Singulières dans les Problèmes aux Limites et en Contrôle Optimal, Lecture Notes in Mathematics # 323, Springer-Verlag, 1973.